The Crystal Structure of α-KTh$_6$F$_{25}$, a Polymorph of CsU$_6$F$_{25}$*

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Crystals of α-KTh$_6$F$_{25}$ are hexagonal-rhombohedral, $R\bar{3}m$, with hexagonal axes $a_0 = 8\cdot313$ (2), $c_0 = 25\cdot262$ (4) Å, $Z = 3$ and ϱ(cal) $= 6\cdot281$ g.cm^{-3}. Structural units of α-KTh$_6$F$_{25}$ are identical to those of CsU$_6$F$_{25}$ and comprise two adjacent rings of Th^{4+}–F polyhedra, 6 polyhedra per ring, with an alkali metal ion centered between the rings. The structural units are stacked $ABCABC$ along c_0 in α-KTh$_6$F$_{25}$ and $ABAB$ in CsU$_6$F$_{25}$. K–F distances are $3\cdot19$ (2) and $3\cdot20$ (3) Å. Th–F distances range from $2\cdot308$ (1) to $2\cdot541$ (9) Å.

Introduction

Compounds exist with the stoichiometry MF.6AF$_4$, where M = K, Rb and Cs, and A = Th and U. Zachariasen (1948) published lattice parameters and the space group ($P6_3/mmc$) for KTh$_6$F$_{25}$ and KU$_6$F$_{25}$, although Asker, Segnit & Wylie (1952) in their work on molten salt system KF–ThF$_4$ and Thoma, Insley, Landau, Friedman & Grimes (1958) in their work on system KF–UF$_4$ did not find an equilibrium compound at KF.6ThF$_4$ or at KF.6UF$_4$. Equilibrium relationships of MF.6AF$_4$ compounds were first established by Thoma *et al.* (1958) in their investigation of the RbF–UF$_4$ system and later by Thoma & Carleton (1961) in their study of system CsF–ThF$_4$. Unpublished work by Thoma and co-workers has established the existence of the remaining complex fluorides (Table 1).

Table 1. *Isomorphous* MF.6AF$_4$ *compounds*

$P6_3/mmc$	$R\bar{3}m$
CsU$_6$F$_{25}$	α-KTh$_6$F$_{25}$
CsTh$_6$F$_{25}$	RbTh$_6$F$_{25}$
RbU$_6$F$_{25}$	
KU$_6$F$_{25}$	
KTh$_6$F$_{25}$	

Brunton (1971) determined the structure of CsU$_6$F$_{25}$ ($P6_3/mmc$), and Table 1 shows the compounds with this structure. The structure of α-KTh$_6$F$_{25}$ has one isomorph, RbTh$_6$F$_{25}$, for which an erroneous space group was given by Brunton, Insley, McVay & Thoma (1965). Although the phase equilibrium work is not complete for all MF–AF$_4$ salt systems, the list in Table 1 is probably exhaustive.

Experimental

The compound α-KTh$_6$F$_{25}$ was grown as primary phase crystals from a molten salt ingot of the composition KF–ThF$_4$ (20–80 mole%). A crystal ground to ellip-

soidal shape (0·156 × 0·156 × 0·182 mm) was mounted on a computer-operated Picker four-circle goniostat, equipped with a scintillation counter detector. Conditions for reflections, hkl, $-h+k+l = 3n$ and the diffraction symmetry $\bar{3}m$, are consistent with space groups $R\bar{3}m$, $R3m$ and $R32$. The solution of the structure was accomplished in $R\bar{3}m$.

Independent reflections were measured by the θ–2θ scan technique using unfiltered Ag $K\alpha$ radiation. Each reflection was step-scanned at intervals of 0·01° with a 2-sec count at each step for a total of 0·5° 2θ. The background was counted for 60 sec on each side of the peak. Take-off angle was 1·5°, and the total range of 2θ for reflections was 3–50°. A standard reflection (060) was measured every 20 reflections to monitor X-ray source variations. The range of net count for (060) was 21855 to 22954.

Lattice parameters were determined by a least-squares adjustment of 12 Ag$^*K\beta$ (0·49701 Å) reflections between 29 and 38° 2θ. At 24°C, $a_o = 8\cdot313$ (2) and $c_o = 25\cdot262$ (4) Å, $Z = 3$ and ϱ(cal) $= 6\cdot281$ g.cm^{-2}. Independent reflections were corrected for Lorentz and polarization effects and absorption ($\mu = 461\cdot2$ cm^{-1}, correction range 0·0156–0·0322).

Starting parameters for Th and K were determined from sections of the three-dimensional Patterson function. When the relationship between the $P6_3/mmc$ and $R\bar{3}m$ structures became evident, atomic parameters for the F ions were estimated from the $P6_3/mmc$ (CsU$_6$F$_{25}$) structure (Table 2). The structure was refined by iterative least-squares adjustment using a modification of the Busing, Martin & Levy (1962) computer program. The quantity minimized by the least-squares program was $\sum w ||F_o^2| - |F_c^2||^2$ with weights, w, equal to the reciprocals of the variances, which were estimated from the empirical equation:

$$\sigma^2(F_o^2) = \{T + B + [0\cdot05(T - B)]^2\}/\{A(\text{Lp})^2\},$$

where T = total counts, B = background, counts A = absorption correction, and Lp = Lorentz and polarization corrections (Brown & Levy, 1964). Scattering factors for the ions were taken from Cromer & Waber (1965), and the anomalous dispersion terms for Th and K for

Ag $K\alpha$ radiation were taken from Cromer (1965). Anisotropic temperature factors were calculated for Th^{4+} and K^+, and the temperature factors for F^- were constrained to be isotropic. Final parameters were calculated using data greater than $\sigma^2(F_o^2)$. Discrepancy indices are:

Table 2. *Atomic parameters of* α-KTh_6F_{25} *and* CsU_6F_{25}

	α-KTh_6F_{25}			CsU_6F_{25}		
	x	y	z	x	y	z
K–Cs	0	0	0	$\frac{1}{3}$	$\frac{2}{3}$	$\frac{3}{4}$
Th–U	0·50031 (7)	$-x$	0·09135 (3)	0·1737	$2x$	0·11146
F(1)	$\frac{1}{2}$	0	0	0·184	$2x$	$\frac{1}{4}$
F(2)	0·193 (2)	$-x$	$-0·0630$ (8)	0·485	$2x$	0·144
F(3)	0·323 (2)	0	$\frac{1}{2}$	0·678	0	0
F(4)	0	0	0·208 (2)	$\frac{1}{3}$	$\frac{2}{3}$	0·068
F(5)	0	0	0·449 (2)	0	0	0·076
F(6)	0·141 (2)	$-x$	$-0·3963$ (9)	0·121	$2x$	0·649

α-KTh_6F_{25} temperature factors

	β_{11}†	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
K	0·001 (2)	β_{11}	0·0001 (2)	$\frac{1}{2}\beta_{11}$	0	0
Th	0·0020 (2)	β_{11}	0·00008 (1)	0·0011 (1)	0·00002 (1)	$-\beta_{13}$
F(1)	0·0053 (3)	Isotropic βs for F; $\beta_{22}=\beta_{11}$,				
F(2)	0·0079 (3)	$\beta_{33}=c^{*2}\beta_{11}/a^{*2}$, $\beta_{12}=b^*\cos\gamma^*\beta_{11}/a^*$,				
F(3)	0·0025 (2)	$\beta_{13}=\beta_{23}=0$. a^*, b^*, c^* and γ^* are the reciprocal lattice parameters.				
F(4)	0·0041 (4)					
F(5)	0·0025 (3)					
F(6)	0·0092 (3)					

† Coefficients in the temperature factor:

$$\exp\left[-(\beta_{11}h^2+\beta_{22}k^2+\beta_{33}l^2+2\beta_{12}hk+2\beta_{13}hl+2\beta_{23}kl)\right]$$

Standard errors in parentheses, corresponding to the last significant digit, are given by the variance–covariance matrix.

Table 3. *Interatomic distances* (Å) *for* α-KTH_6F_{25}

6[K——F(6)]	3·19 (2)	6[K——F(2)]	3·20 (3)
Th——F(1)	2·308 (1)	2[Th—F(3)]	2·330 (4)
2[Th—F(2)]	2·34 (1)	2[Th—F(6)]	2·35 (1)
Th——F(4)	2·481 (7)	Th——F(5)	2·541 (9)
4[F(1)-F(2)]	2·74 (2)	4[F(1)-F(6)]	2·75 (2)
F(2)——F(4)	2·56 (3)	2[F(2)-F(6)]	2·77 (2)
2[F(3)-F(3)]	2·69 (2)	[F(5)—F(6)]	2·43 (3)
F(5)——F(5)	2·58 (6)		

Comparable cation–anion distances (Å) for CsU_6F_{25}

6[Cs—F(2)]	3·12 (2)	6[Cs——F(6)]	3·45 (4)
2[U——F(3)]	2·267 (8)	U———F(1)	2·279 (2)
2[U——F(6)]	2·29 (5)	2[U—F(2)]	2·29 (2)
U———F(4)	2·39 (4)*	U———F(5)	2·54 (6)

*Accidentally omitted from Table 2 (Brunton, 1971).
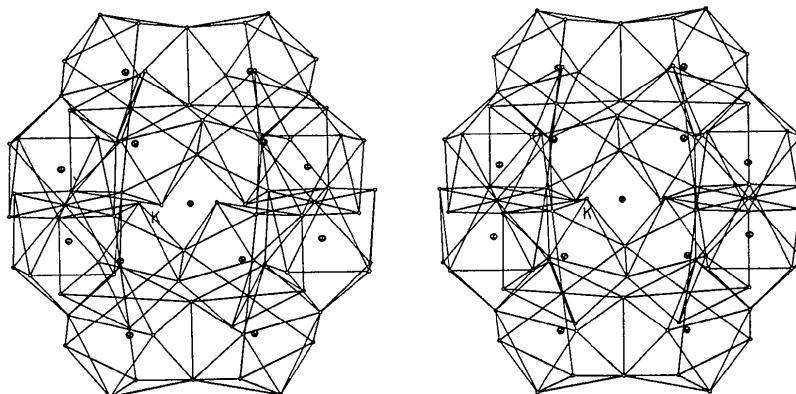
Fig. 1. One double ring of Th–F polyhedra with a central K ion. Ring tilted 30° about Cartesian y axis.

$R_1 = \sum ||F_o^2| - F_c^2||/\sum|F_o^2|$ $= 0\cdot0802$ for reflections $> \sigma$

                           $= 0\cdot0834$ for all 736 reflections.

$R_2 = \sum ||F_o| - |F_c||/\sum|F_o|$ $= 0\cdot0616$ for 641 reflections $> \sigma$

                           $= 0\cdot0782$ for all reflections.

The standard deviations of an observation of unit weight $[\sum w(F_o - F_c)^2/(n_o - n_v)]^{1/2}$ is $1\cdot6024$ where n_o is the number of reflections, 641, and n_v is the number of variables. Atomic parameters and temperature factors are listed in Table 2. Interatomic distances are given

Table 4. Observed and calculated structure factors for α-KTh$_6$F$_{25}$

[Dense multi-column table of structure factors with columns L, FOBS, FCAL repeated]

in Table 3. Observed and calculated structure factors are listed in Table 4. An extinction correction was applied to F_c by the method suggested by Zachariasen (1967, 1968); Zachariasen's $r^* = 0\cdot00939$ (8).

Results

The basic structural unit of α-5Th$_6$F$_{25}$ is a double ring of Th–F polyhedra, six polyhedra per ring, surrounding the K ion (Fig. 1). Each Th–F polyhedron is a capped trigonal prism with an F ion at each of the nine corners. Rings are joined by opposing polyhedra, sharing corners, F(1), as with CsU$_6$F$_{25}$ (Brunton, 1971). The result is that each K is surrounded by 12F at distances of $3\cdot19$ (2) and $3\cdot20$ (3) Å. Within each ring the polyhedra link by sharing edges F(2)–F(4) and F(6)–F(5), alternately, and these are short F–F distances. Interatomic distances for both compounds are listed in Table 3 for comparison and to correct an omission in the original CsU$_6$F$_{25}$ table. Th–F distances, except for Th–F(5) [$2\cdot541$ (9) Å] are 0·03 to 0·1 Å longer than corresponding U–F distances. Cs–F distances are slightly longer than corresponding K–F distances.

The difference in the structures of α-KTh$_6$F$_{25}$ and CsU$_6$F$_{25}$ results from a different stacking arrangement for the double ring in each compound. The stacking order for CsU$_6$F$_{25}$ along c_o is $ABAB$ [$c_o = 16\cdot78$ (2) Å] and the stacking order for α-KTh$_6$U$_{25}$ is $ABCABC$ $c_o = 25\cdot262$ (2) $\simeq \frac{3}{2}$ (16·78) Å]. The stacking order should not be confused with hexagonal or cubic close-packing, because atoms within the structural units are not close-packed. Structures are illustrated in Fig. 2 & 3; the F polyhedra are omitted for clarity. Preliminary work by Thoma *et al.* (unpublished) indicates that the difference in the structures represents a polymorphic transition in the compound KTh$_6$F$_{25}$, and that the KTh$_6$F$_{25}$ reported by Zachariasen (1948) is the low temperature or β-KTh$_6$F$_{25}$ (CsU$_6$F$_{25}$ form).

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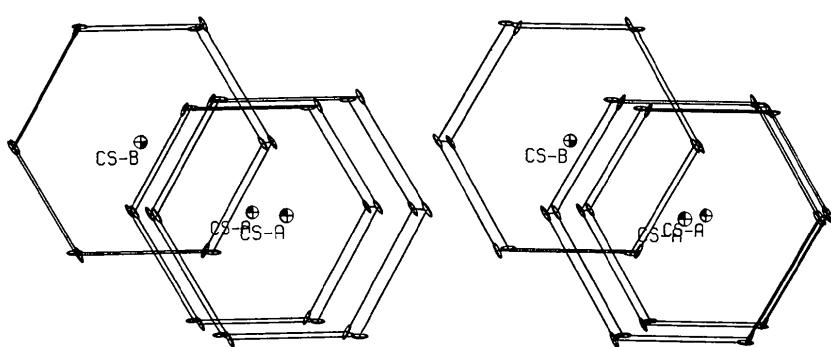
Fig. 2. *ABA* stacking of CsU$_6$F$_{25}$. F ions are omitted. U ions are not labeled.
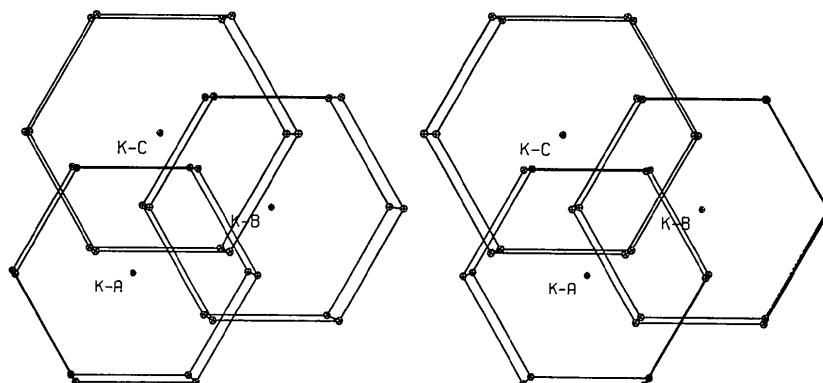
Fig. 3. *ABC* stacking of α-KTh_6F_{25}. F ions are omitted. Th ions are not labeled.

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The Crystal Structure of the Molecular Complex between Antimony Trichloride and Phenanthrene

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$2SbCl_3$.phenanthrene gives triclinic ($P\bar{1}$) crystals: $a=8\cdot990(6)$, $b=11\cdot10(1)$, $c=10\cdot44(1)$ Å, $\alpha=90\cdot7(1)$, $\beta=80\cdot4(1)$, $\gamma=69\cdot1(1)°$, $Z=2$. The structure has been determined by three-dimensional Fourier methods and refined using block-diagonal least-squares down to $R=3\cdot8\%$. The structure consists of alternating double layers of $SbCl_3$ and phenanthrene molecules. The coordination around Sb corresponds to a distorted trigonal bipyramid. Two independent Sb atoms are on opposite sides of the phenanthrene plane and their distances from that plane are not equal ($3\cdot27$, $2\cdot94$ Å). The interactions between the metal atom and the aromatic molecule concern the π-electron system. Distorted octahedral coordination around the Sb atoms is completed by Sb$\cdots$Cl contacts ($3\cdot260$, $3\cdot49$, $3\cdot41$, $3\cdot55$ Å) significantly shorter than the sum of van der Waals radii.

Introduction

It is of interest to study the structure of the complexes formed by antimony trihalides with aromatic hydrocarbons, to determine the nature of the bonding interactions between the constituents and to obtain information on the mechanism of the catalytic effect which the antimony trihalides exert in Friedel–Crafts type reactions. The structure study of $2SbCl_3$.phenanthrene has been carried out to make a contribution to

understanding the nature of the interactions between the inorganic and aromatic constituents, which could be related to the donor properties of the π system of phenanthrene and to the acceptor properties of $SbCl_3$. From the study of the pure quadrupole resonance spectrum of $2SbCl_3$.benzene, Okuda, Nakao, Shiroyama & Negita (1968) suggested that in this compound there is a charge transfer from the π system of benzene to a chlorine atom of $SbCl_3$, but a subsequent X-ray analysis carried out by Hulme & Szymanski (1969) showed that